Three-dimensional analytical periodic solutions of the Laplace equation

This article has been downloaded from IOPscience. Please scroll down to see the full text article.
1985 J. Phys. A: Math. Gen. 18 L845
(http://iopscience.iop.org/0305-4470/18/14/004)

View the table of contents for this issue, or go to the journal homepage for more

Download details:
IP Address: 129.252.86.83
The article was downloaded on 31/05/2010 at 09:01

Please note that terms and conditions apply.

## LETTER TO THE EDITOR

# Three-dimensional analytical periodic solutions of the Laplace equation 

D Ouroushev<br>University of Sofia, Faculty of Physics, Department of Solid State Physics, Boulevard A Ivanov 5, Sofia 1126, Bulgaria

Received 17 May 1985


#### Abstract

A method is proposed, by which three-dimensional periodic analytical solutions of the Laplace equation can be found. The solutions obtained describe the electrostatic potential in a three-dimensional space lattice of point charges with a certain symmetry.


The problem of finding the three-dimensional solution of the Laplace equation

$$
\begin{equation*}
\Delta \psi=0 \tag{1}
\end{equation*}
$$

is usually solved by separating the variables. By this method a general solution can be obtained, which has the form (Morse et al 1953)

$$
\begin{equation*}
\psi=\mathrm{e}^{ \pm i r x} \mathrm{e}^{ \pm i s y} \mathrm{e}^{ \pm t z}, \quad t^{2}=r^{2}+s^{2} \tag{2}
\end{equation*}
$$

As can be seen from (2) this solution is periodic in the $x$ and $y$ directions and exponentially decreasing or increasing in the $z$ direction. It should be mentioned here that the correlation between the constants $r, s$ and $t$ is always such that the coefficient of $z$ is real. Consequently the method of separation of the variables cannot be used to obtain a solution periodic in three dimensions.

Here a method will be proposed by which a three-dimensional periodic solution of the Laplace equation can be found. Let us make the following substitution in equation (1)

$$
\begin{equation*}
\psi=4 q \operatorname{Arth}\left[u_{1}(x) v_{1}(y) w_{1}(z)+u_{2}(x) v_{2}(y) w_{2}(z)\right] \tag{3}
\end{equation*}
$$

where $u_{i}, v_{i}, w_{i}(i=1,2)$ are Jacobi elliptic functions which satisfy the following nonlinear ordinary differential equations (Janke et al 1960)

$$
\begin{align*}
& \left(\mathrm{d} u_{i} / \mathrm{d} x\right)^{2}=A_{i}^{x} u_{i}^{4}+B_{i}^{x} u_{i}^{2}+C_{i}^{x} \\
& \left(\mathrm{~d} v_{i} / \mathrm{d} y\right)^{2}=A_{i}^{y} v_{i}^{4}+B_{i}^{y} v_{i}^{2}+C_{i}^{y} \quad i=1,2  \tag{4}\\
& \left(\mathrm{~d} w_{i} / \mathrm{d} z\right)^{2}=A_{i}^{z} w_{i}^{4}+B_{i}^{2} w_{i}^{2}+C_{i}^{z} .
\end{align*}
$$

Substituting (4) into equation (1) we obtain

$$
\begin{aligned}
\left(A_{1}^{x} u_{1}^{2}+A_{1}^{y} v_{1}^{2}\right. & \left.+A_{1}^{z} w_{1}^{2}\right) 2 \alpha+\left(A_{2}^{x} u_{2}^{2}+A_{2}^{y} v_{2}^{2}+A_{2}^{z} w_{2}^{2}\right) 2 \beta \\
& -2\left(A_{1}^{x} u_{1}^{2}+A_{1}^{y} v_{1}^{2}+A_{1}^{z} w_{1}^{2}\right)^{2}\left(\alpha \beta^{2}+\beta^{2} \alpha\right)+\left(B_{1}^{x}+B_{1}^{y}+B_{1}^{2}\right)\left(\alpha+\alpha^{3}+\alpha \beta^{2}\right) \\
& +\left(B_{2}^{x}+B_{2}^{y}+B_{2}^{z}\right)\left(\beta+\beta^{3}+\beta^{2} \alpha\right)+2(\alpha+\beta)
\end{aligned}
$$

$$
\begin{align*}
& \times\left(C_{1}^{x} v_{1}^{2} w_{1}^{2}+C_{2}^{x} v_{2}^{2} w_{2}^{2}+C_{1}^{y} u_{1}^{2} w_{1}^{2}+C_{2}^{y} u_{2}^{2} w_{2}^{2}+C_{1}^{2} u_{1}^{2} v_{1}^{2}+C_{2}^{z} u_{2}^{2} v_{2}^{2}\right) \\
& +4(\alpha+\beta)\left(u_{1}^{\prime} u_{2}^{\prime} v_{1} w_{1} v_{2} w_{2}+v_{1}^{\prime} v_{2}^{\prime} u_{1} u_{2} w_{1} w_{2}+w_{1}^{\prime} w_{2}^{\prime} u_{1} u_{2} v_{1} v_{2}\right)=0 . \tag{5}
\end{align*}
$$

Here the prime denotes a derivative by the corresponding variable; $\alpha$ and $\beta$ are correspondingly

$$
\begin{equation*}
\alpha=u_{1} v_{1} w_{1} \quad \beta=u_{2} v_{2} w_{2} . \tag{6}
\end{equation*}
$$

Equation (5) can be further simplified by using some concrete properties of the Jacobi elliptic functions. Let us assume that $u_{i} v_{i} w_{i}$ are some of the three main Jacobi elliptic functions $\operatorname{sn}(x, k), \operatorname{cn}(x, k), \operatorname{dn}(x, k)$ (Janke et al 1960). In this case

$$
\begin{align*}
& u_{1}^{\prime} u_{2}^{\prime}=u_{1} u_{2}\left(D_{1}^{x}+D_{2}^{x}\right)+A_{1}^{x} u_{1}^{2} u_{1} u_{2}+A_{2}^{x} u_{2}^{2} u_{1} u_{2} \\
& v_{1}^{\prime} v_{2}^{\prime}=v_{1} v_{2}\left(D_{1}^{\prime}+D_{2}^{y}\right)+A_{1}^{y} v_{1}^{2} v_{1} v_{2}+A_{2}^{v} v_{2}^{2} v_{1} v_{2}  \tag{7}\\
& w_{1}^{\prime} w_{2}^{\prime}=w_{1} w_{2}\left(D_{1}^{z}+D_{2}^{z}\right)+A_{1}^{z} w_{1}^{2} w_{1} w_{2}+A_{2}^{z} w_{2}^{2} w_{1} w_{2} .
\end{align*}
$$

Of course the relations (7) are valid if the pairs of functions $u_{1} u_{2} ; v_{1} v_{2} ; w_{1} w_{2}$ depend on the same modal $k_{x}, k_{y y} k_{z}$ respectively. In (7) $D_{1}^{x y z}, D_{2}^{x y z}$ are constants, which can be determined by making a concrete choice for the functions $u_{i} v_{i} w_{i}$; they can be expressed by the coefficients $B_{i}^{x y z}$.

It must be underlined that for every combination of main elliptic functions ( sncn ; $\mathrm{sn} \mathrm{dn} ; \mathrm{cn} \mathrm{dn}$ ) relations (7) are valid.

Using (7) equation (5) can be simplified:

$$
\begin{align*}
\left(B_{1}^{x}+B_{1}^{y}+B_{1}^{z}\right) & \left(\alpha+\alpha^{3}\right)+\left(B_{2}^{x}+B_{2}^{y}+B_{2}^{z}\right)\left(\beta+\beta^{3}\right) \\
& +\left[2\left(D_{1}^{x}+D_{2}^{x}\right)+2\left(D_{1}^{y}+D_{2}^{y}\right)+2\left(D_{1}^{z}+D_{2}^{z}\right)-B_{1}^{x}-B_{1}^{y}-B_{1}^{z}\right] \alpha \beta^{2} \\
& +\left[2\left(D_{1}^{x}+D_{1}^{y}\right)+2\left(D_{1}^{y}+D_{2}^{y}\right)+2\left(D_{1}^{z}+D_{2}^{z}\right)-B_{2}^{x}-B_{2}^{y}-B_{2}^{z}\right] \beta^{2} \alpha \\
& +\left[A_{1}^{x} u_{1}^{2}+A_{1}^{y} v_{1}^{2}+A_{1}^{z} w_{1}^{2}+C_{1}^{x} v_{1}^{2} w_{1}^{2}+C_{2}^{x} v_{2}^{2} w_{2}^{2}+C_{1}^{y} u_{1}^{2} w_{1}^{2}\right. \\
& \left.+C_{2}^{y} u_{2}^{2} w_{2}^{2}+C_{1}^{z} u_{1}^{2} v_{1}^{2}+C_{2}^{z} u_{2}^{2} w_{2}^{2}\right] 2 \alpha+\left[A_{2}^{x} u_{2}^{2}+A_{2}^{y} v_{2}^{2}\right. \\
& +A_{2}^{z} w_{2}^{2}+C_{1}^{x} v_{1}^{2} w_{1}^{2}+C_{2}^{x} v_{2}^{2} w_{2}^{2}+C_{1}^{y} u_{1}^{2} w_{1}^{2}+C_{2}^{y} u_{2}^{2} w_{2}^{2} \\
& \left.+C_{1}^{z} u_{1}^{2} v_{1}^{2}+C_{2}^{z} u_{2}^{2} v_{2}^{2}\right] 2 \beta=0 . \tag{8}
\end{align*}
$$

Consequently the problem of solving the three-dimensional Laplace equation is reduced to that of solving a system of algebraic equations. The first three equations of this system, as follows from (8), are

$$
\begin{align*}
& B_{1}^{x}+B_{1}^{y}+B_{1}^{z}=0 \\
& B_{2}^{x}+B_{2}^{y}+B_{2}^{z}=0  \tag{9}\\
& D_{1}^{x}+D_{2}^{x}+D_{1}^{y}+D_{2}^{y}+D_{1}^{z}+D_{2}^{z}=0 .
\end{align*}
$$

Moreover, the coefficients of $2 \alpha$ and $2 \beta$ in (8) must be equal to zero:

$$
\begin{align*}
A_{1}^{x} u_{1}^{2}+A_{1}^{y} v_{1}^{2}+ & A_{1}^{z} w_{1}^{2}+C_{1}^{x} v_{1}^{2} w_{1}^{2}+C_{2}^{x} v_{2}^{2} w_{2}^{2} \\
& +C_{1}^{y} u_{1}^{2} w_{1}^{2}+C_{2}^{y} u_{2}^{2} w_{2}^{2}+C_{1}^{z} u_{1}^{2} v_{1}^{2}+C_{2}^{z} u_{2}^{2} v_{2}^{2}=0 \\
A_{2}^{x} u_{2}^{2}+A_{2}^{y} v_{2}^{2}+ & A_{2}^{z} w_{2}^{2}+C_{1}^{x} v_{1}^{2} w_{1}^{2}+C_{2}^{x} v_{2}^{2} w_{2}^{2}  \tag{10}\\
& +C_{1}^{y} u_{1}^{2} w_{1}^{2}+C_{2}^{y} u_{2}^{2} w_{2}^{2}+C_{1}^{z} u_{1}^{2} v_{1}^{2}+C_{2}^{z} u_{2}^{2} v_{2}^{2}=0 .
\end{align*}
$$

Equations (10) lead to a further ten relations between the coefficients $A_{i}^{x} A_{i}^{y} A_{i}^{z} C_{i}^{x} C_{i}^{y} C_{i}^{z}$ ( $i=1,2$ ). The exact form of these equations can be given after a concrete choice of the functions $u_{i} v_{i} w_{i}$. Solving the system of algejraic equations determined by ( 9 ) and (10) we can in principle find a three-dimensional periodic solution of the Laplace equation, expressed in Jacobi elliptic functions.

Let us now do a concrete choice for the functions $u_{i} v_{i} w_{i}$ setting

$$
\begin{array}{ll}
u_{1}=a_{1}^{x} \operatorname{cn}\left(l x, k_{x}\right) & u_{2}=a_{2}^{x} \operatorname{sn}\left(l x, k_{x}\right) \\
v_{1}=a_{1}^{y} \operatorname{cn}\left(m y, k_{y}\right) & v_{2}=a_{2}^{y} \operatorname{sn}\left(m y, k_{y}\right)  \tag{11}\\
w_{1}=a_{1}^{2} \operatorname{cn}\left(n z, k_{z}\right) & w_{2}=a_{2}^{2} \operatorname{dn}\left(n z, k_{z}\right) .
\end{array}
$$

As can be seen from (11) 18 indeterminated constants $A_{i}^{x y z} B_{i}^{x y z} C_{i}^{x y z}$ reduce to 12 due to the fact, that for every pair of functions $u_{i}, v_{i}, w_{i}(i=1,2)$ the coefficients before $x y z$, as the modals $k_{x}, k_{y}, k_{z}$ must be the same, because in (9) and (10) these functions must be combinate.

Due to the fact that the function $u_{i} v_{i} w_{i}$ occurs in the solution only in the combinations $u_{1} v_{1} w_{1}, u_{2} v_{2} w_{2}$, the coefficients $a_{i}^{x}, a_{i}^{y}, a_{i}^{z}, a_{2}^{x}, a_{2}^{y}, a_{2}^{z}$ can be reduced to two:

$$
\begin{equation*}
A=a_{1}^{x} a_{1}^{y} a_{1}^{z} \quad B=a_{2}^{x} a_{2}^{y} a_{2}^{z} \tag{12}
\end{equation*}
$$

In this case the system of equations (9) and (10) reduces to eight equations for the eight coefficients $A, B, k_{x}, k_{y}, k_{z}, l, m, n$. Solving this system we obtain

$$
\begin{equation*}
A=B=1 \quad l^{2}+m^{2}=n^{2} \quad k_{x}^{2}=k_{y}^{2}=1-k_{z}^{2} . \tag{13}
\end{equation*}
$$

Consequently the function
$\psi=4 q \operatorname{Arth}\left[\mathrm{cn}(l x, k) \mathrm{cn}(m y, k) \mathrm{cn}\left(n z, k^{\prime}\right)+\mathrm{sn}(l x, k) \operatorname{sn}(m y, k) \operatorname{dn}\left(n z, k^{\prime}\right)\right]$
where $k^{2}=1-k^{\prime 2}$, is a three-dimensional periodic solution of the Laplace equation. The solution obtained is periodic in the $x, y$ and $z$ direction with respective periods

$$
\begin{equation*}
T_{x}=4 K(k) / l \quad T_{y}=4 K(k) / m \quad T_{z}=4 K\left(k^{\prime}\right) / n \tag{15}
\end{equation*}
$$

where $K(k)$ and $K\left(k^{\prime}\right)$ are the full elliptic integrals of the first kind determining the periods of the elliptic functions.

It must be mentioned that the obtained function (14) possesses singularities at the points in which the argument of the function Arth is equal to $\pm 1$. These singularities can be interpreted according to the physical meaning of the Laplace equation. This equation describes the electrostatic potential in a system of charges in the areas where the space charge density is zero.

The presence of singularities in the solution of the Laplace equation is usually connected with the existence of a point charge in these points (Jackson 1962). Setting the constant $q$ in (3) equal to the absolute value of these point charges it can be said that a solution of this type describes the electrostatic potential in a system of point charges distributed periodically and forming a space lattice.

For the concrete solution (14) this space lattice is given in figure 1. As can be seen from the figure, this is a space lattice from the rhomboid system with $C$ primitive cell.

Consequently the solution obtained is an analytical expression for the electrostatic potential in the crystal-like structure of point charges with the aforementioned symmetry.

Here it must be mentioned that the proposed method gives us a possibility of finding solutions with other symmetry, which are also expressed in Jacobi elliptic


Figure 1. The space lattice of point charges described by the solution (14) of the Laplace equation. Open circles, positive point charges; full circles, negative point charges.
functions. Consequently the proposed method makes it possible to find an analytical expression for the electrostatic potential in a three-dimensional periodic structure of point charges or we obtain an analytical expression for the crystal field in a certain type of ionic crystal with corresponding symmetry.

## References

Jackson J D 1962 Classical electrodynamics (New York: Wiley)
Janke E, Emde F and Lösch F 1960 Tafeln Höherer Funktionen (Stuttgart: Teubuer)
Morse P H and Feshbach H 1953 Methods of Theoretical Physics Part II (New York: McGraw-Hill)

